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## Study of $W_4$ - curvature tensor in the space-time of general relativity

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#### Abstract

In this paper, we study the physical properties of  $W_4$ - Curvature tensor and found that it does not satisfy the Bianchi differential identity even with Ricci tensor being of Codazzi type. It is further shown that Einstein field equations with Cosmological constant can be expressed with the help of  $W_{4\beta\gamma}$  and Rainich conditions for the existence of the non-null electrovariance can also be written with the help of  $W_{4ij}$ . Finally, it is shown that  $W_4$ -flat space time signifies empty field.

**Keywords:** W<sub>4</sub> curvature tensor, Einstein field equations, Rainich conditions, W<sub>4</sub>-flat space time. AMS 2021 subject classification: 53C50, 53C80

#### Introduction

On the line Wyel's of projective curvature tensor Pokhariyal (1973) [1] defined  $W_3$  and  $W_4$  curvature tensors, to study geometric and physical properties of these tensors in the 4-dimensional space time of general relativity.

$$W_3(X,Y,Z,T) = R(X,Y,Z,T) + \frac{1}{n-1} [g(Y,Z)Ric(X,T) - g(Y,T)Ric(X,Z)]$$
 (1.1)

$$W_4(X,Y,Z,T) = R(X,Y,Z,T) + \frac{1}{n-1} [g(X,Z)Ric(Y,T) - g(X,Y)Ric(T,Z)]$$
 (1.2)

Moindi (2007) [2] have studied the geometric properties of  $W_3$  tensor in Lorentzian Para-Sasakian and other complex manifolds. The  $W_4$ -curvature tensor, has been studied in Sasakian and LP-Sasakian manifolds (Katende, 2011) [3]. Additional physical properties of  $W_4$ -curvature tensor in the space time of general relativity are studied in this paper.

#### Properties of $W_4$ -curvature tensor

The  $W_4$ -curvature tensor has been defined by (1.2) where R(X,Y,Z,T) = g(R(X,Y,Z),T) and  $R(X,Y,Z) = D_X D_Y Z - D_Y D_X Z - D_{[X,Y]} Z$  is the Riemann curvature tensor, R(X,Y) = g(R(X),Y) is the (0,2)-type Ricci tensor, R is the scalar curvature and R denotes the Riemannian connection. It is seen Pokhariyal, 1973 [1], that  $W_4(X,Y,Z,T) \neq W_4(Y,X,Z,T)$ ,  $W_4(X,Y,Z,T) \neq W_4(X,Y,Z,T) \neq W_4(X,Y,Z,T)$ 

Therefore, none of symmetric or skew-symmetric properties are satisfied by  $W_4$ -curvature tensor.

From (1.2), it is seen that

$$W_4(X,Y,Z,T) + W_4(X,Z,T,Y) + W_4(X,T,Y,Z) = 0. (2.1)a$$

$$W_4(X,Y,Z,T) + W_4(Y,Z,X,T) + W_4(Z,X,Y,T)0$$
(2.1)b

That is, both the cyclic properties are satisfied by  $W_4$ -curvature tensor.

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#### **Bianchi Differential Identity**

It is shown that the Bianchi differential identity is given by

$$(\nabla_{U}R)(X,Y,Z,T) + (\nabla_{Z}R)(X,Y,T,U) + (\nabla_{T}R)(X,Y,U,Z) = 0$$
(3.1)

For the four-dimensional space time of general relativity, from (1.2), we have

$$W_4(X,Y,Z,T) = R(X,Y,Z,T) + \frac{1}{3} [g(X,Z)Ric(Y,T) - g(X,Y)Ric(Z,T)]$$
(3.2)

In order to check for Bianchi differential identity for  $W_4$ -curvature tensor, we compute

$$(\nabla_{U}W_{4})(X,Y,Z,T) + (\nabla_{Z}W_{4})(X,Y,T,U) + (\nabla_{T}W_{4})(X,Y,U,Z) = (\nabla_{U}R)(X,Y,Z,T) + \frac{1}{3} [g(X,Z)(\nabla_{U}Ric)(Y,T) - g(X,Y)(\nabla_{U}Ric)(Z,T)] + (\nabla_{Z}R)(X,Y,T,U) + \frac{1}{3} [g(X,T)(\nabla_{Z}Ric)(Y,U) - g(X,Y)(\nabla_{Z}Ric)(T,U)] + (\nabla_{T}R)(X,Y,U,Z) + \frac{1}{3} [g(X,U)(\nabla_{T}Ric)(Y,Z) - g(X,Y)(\nabla_{T}Ric)(U,Z)]$$
(3.3)

By virtue of (3.1) and (3.2) we have.

$$(\nabla_{U}W_{4})(X,Y,Z,T) + (\nabla_{Z}W_{4})(X,Y,T,U) + (\nabla_{T}W_{4})(X,Y,U,Z) = \frac{1}{3}[g(X,Z)(\nabla_{U}Ric)(Y,T) + g(X,T)(\nabla_{Z}Ric)(Y,U) + g(X,U)(\nabla_{T}Ric)(Y,Z) - g(X,Y)\{(\nabla_{U}Ric)(Z,T) + (\nabla_{Z}Ric)(U,T) + (\nabla_{T}Ric)(U,Z)\}$$
(3.4)

The condition for Ricci tensor to be of Codazzi type, Derdzinski and Shen, 1983 [4] is given by

$$(\nabla x \operatorname{Ric})(Y, Z) = (\nabla y \operatorname{Ric})(X, Z) = (\nabla z \operatorname{Ric})(X, Y) \tag{3.5}$$

It is noticed that using (3.5), the equation (3.4) does not vanish. Hence,  $W_4(X, Y, Z, T)$  does not satisfy the Bianchi differential identity, even with Ricci tensor being Codazzi type.

#### **Einstein Field Equations**

The Einstein field equations with cosmological term is given by

$$R_{\beta}\gamma - \frac{1}{2}Rg_{\beta}\gamma + \Lambda g_{\beta}\gamma = kT_{\beta},\tag{3.6}$$

where  $\Lambda$  is the cosmological constant, K is the non-zero gravitational constant and  $T_{\beta}\gamma$  is the energy momentum tensor. Writing  $W_4$  tensor in the local coordinates, we have from (3.2)

$$W_{4\alpha\beta\gamma\delta} = R_{\alpha\beta\gamma\delta} + \frac{1}{3} [g_{\alpha\gamma}R_{\beta\delta} - g_{\alpha\beta}R_{\gamma\delta}]. \tag{3.7}$$

$$W_{4\beta\gamma\delta}^{\alpha} = R_{4\beta\gamma\delta}^{\alpha} + \frac{1}{3} [g_{\gamma}^{\alpha} R_{\beta\delta} - g_{\beta}^{\alpha} R_{\gamma\delta}] \tag{3.8}$$

Contracting for  $\alpha$  and  $\delta$ , we get

$$W_{4\beta\gamma} = R_{\beta\gamma} + \frac{1}{3} [R_{\beta\gamma} - R_{\beta\gamma}] = R_{\beta\gamma} \tag{3.9}$$

Further contraction gives

$$W_4 = R \tag{3.10}$$

Hence, the Einstein tensor and the Einstein field equations with cosmological constant can alternatively be expressed using (3.9) and (3.10) with the help of  $W_{4\beta\gamma}$  and  $W_4$  as follows:

$$E_{\beta\gamma} = W_{4\beta\gamma} - \frac{1}{2}W_4g_{\beta\gamma} \tag{3.11}$$

and

$$W_{4\beta\gamma} - \frac{1}{2}W_4g_{\beta\gamma} + \Lambda g_{\beta\gamma} = kT_{\beta\gamma} \tag{3.12}$$

#### W<sub>4</sub> -Flat Space Time

We now study  $W_4$  flat Space Time and some other physical properties.

**Definition:** A space time is said to be  $W_4 - flat$  if tensor  $W_4$  vanishes in it.

Thus, for a  $W_4$  flat Space Time we have from (3.8)

$$R^{\alpha}_{\beta\gamma\delta} = \frac{1}{2} [g^{\alpha}_{\beta} R_{\gamma\delta} - g^{\alpha}_{\gamma} R_{\beta\delta}] \tag{3.13}$$

On contracting  $\alpha$  and  $\delta$ , we get

$$R_{\beta\gamma} = \frac{1}{3} [R_{\beta\gamma} - R_{\beta\gamma}] = 0 \tag{3.14}$$

Thus, a  $W_4$  flat space time signifies empty field, which may contain isolated masses as singularities. It can also be said that empty space can still be curved and gravity can still be felt.

The vector

$$\theta_i = \frac{g_{ij} \in j^{klm} R_k^p R_{pl;m}}{\sqrt{-g} R_{ab} R^{ab}},\tag{3.15}$$

known as complexion vector of a non-null electromagnetic field with no matter (Misner and Wheeler, 1954 [5] and its vanishing implies that the field is purely electrical. It was shown [1] that the vanishing of the divergence of  $W_{4ijk}^h$  in an electromagnetic field implies purely electric field.

Rainich (1952) [6] has shown that the necessary and sufficient condition for the existence of the non-null electro variance are

$$R = 0 ag{3.16}$$

$$R_j^i R_k^j = \frac{1}{4} \delta_k^i R_{ab} R^{ab} \tag{3.17}$$

$$\theta_{i;j} = \theta_{j;i} \tag{3.18}$$

From (3.9) and (3.10), it is noticed that  $R_{ij}$  can very well be replaced by  $W_{4ij}$  and R be replaced by  $W_4$  in expressing the Rainich conditions.

#### Discussion

One of the advantages of  $W_4$ -Curvature tensor is that its contracted part is equal to Ricci tensor and the scalar  $W_4$  is equal to R, the scalar curvature. Therefore, these results can be alternatively used for getting various geometrical and physical results. Further, the  $W_4$ -flat space time signifies empty field.  $W_4$ -Curvature tensor has additional terms as compared to Riemann curvature tensor and is likely to provide additional interpretations of geometrical and physical phenomena.

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