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Mathematical model on two phase of hepatic blood flow in venules with special reference to malaria

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Abstract

In this paper, we analyzed blood flow in hepatic venules. The viscosity of blood increases in venules due to more formation of rouleaux along the axis of red blood cells as at the same venules are remote from the heart and proximate to liver so, we applied the Herschel-Bulkley non-Newtonian model in Bio-fluid mechanical setup. Employing Navier-Stoke equation and equation of continuity in cylindrical co-ordinate system and all required mathematical formations are in tensorial form. Using by numerical method we calculated the value of parameter for a clinical data. Finally, we obtained a linear relationship between blood pressure drop and hematocrit for particular value of parameter and discussed in the graph.

Keywords: Hepatic circulation, malaria, hematocrit, viscosity, venule, Rouleaux

1. Introduction

1.1 Structure and Function of Liver

The liver has the most complicated circulation of any one of the organ. According to the anatomical peculiarity of the double afferent blood supply of the liver, 75%-80% of the blood entering the liver is partially deoxygenated venous blood provided by the portal vein, which collects all the blood that leaves the spleen, stomach, small and large intestine, gallbladder and pancreas (Vollmar, B. *et al.* 2009, Rappaport, AM *et al.* 1980)^[24].

The liver has a wide range of functions including detoxification of various metabolites and toxic matter regulation of glycogen storage, decomposition of red blood cells, hormone production and the production of biochemicals necessary for digestion and other metabolic activities. The liver is the only human internal organ capable of natural regeneration of lost tissues as little as 25% of a liver can regenerate into a whole liver ^[10].

The portal venous system is responsible for directing blood from parts of the gastrointestinal tract of the liver. Blood flow to the liver is unique in that it receives both oxygenated and deoxygenated blood. Blood passes from branches of the portal vein through cavities between "plates" of hepatocytes called sinusoids blood also flows from branches of the hepatic artery and mixes in the sinusoids to supply the hepatocyte with oxygen. This mixture percolates through the sinusoids and collects in a central vein which drains into the hepatic vein. The hepatic vein subsequently drains into the inferior vena cava. The hepatic artery provides 30 to 40% of the oxygen to the liver, while only accounting for 25% of the total liver blood flow. The rest comes from the partially deoxygenated blood from the portal vein ^[3, 4].

1.2 Structure and Function of Venule

The hepatic microcirculation generally refers to the circulatory system beginning with portal venule, extending to the terminal portal venule, and then reaching the sinusoid network, followed by the postcapillary terminal hepatic venule and ending with the muscular venule.^[27] The basic structural unit of the hepatic microcirculation is the hepatic lobule, in which a terminal hepatic venule is located at the center and several portal venules at the periphery, with the hepatic sinusoids running from the terminal portal venule, forming a hexagonal vascular structure ^[28].

1.3 Composition of Blood

Blood is bio fluid or fluid connective tissue. Blood consists of a suspension of cells in an aqueous solution called plasma which composed of about 90% water and 7% protein. There are about 95% are red cells or erythrocytes whose main function is to transport oxygen from lungs to all the cells of the body and removal of carbon dioxide formed by metabolic process in the body to lungs. About 45% of the blood volume in an average human is occupied by red cells. This fraction is known as the hematocrit of the remaining white cells or leucocytes constitute about one sixth or 1% of total and these play an impartment role in the resistance of the body to infection and platelets form 5% of the total blood and they perform a function related to blood clotting ^[2, 18].

1.4 Description of Disease and RBC deformability in Malaria

Malaria is caused by single celled microorganisms of the *Plasmodium* group. It is spread exclusively through bites of infected *Anopheles* mosquitoes. The mosquito bite introduces the parasites from the mosquito's saliva into a person's blood. The parasites travel to the liver where they mature and reproduce ^[31]. The Plasmodium life cycle is very complex and takes place in two phases; sexual and asexual, the vector mosquitoes and the vertebrate hosts. In the vectors, mosquitoes, the sexual phase of parasite's life cycle occurs. The asexual phase of life cycle occurs in humans, the intermediate host for malaria ^[29, 30].

Red blood cells host plasmodium parasites that cause malaria, of which plasmodium falciparum is the most pathogenic. The deformability of RBC is markedly modified by invasion and development of *P. falciparum*. The deformability of RBC depends on three parameters: (i) the membrane elasticity (ii) the cytoplasmic viscosity that depends on intracellular ion and haemoglobin concentration and (iii) the surface to volume ratio. The balance among these three parameters can be altered during malaria (Lavazec, 2017).

2. Real Model

2.1 Frame of reference

In this model, we employ Navier-Stoke's equation and equation of continuity and have chosen orthogonal curvilinear generalized three-dimensional coordinate system denoted by E³ called three dimensional Euclidean space of the moving blood. All quantities related to blood flow are written in tensorial form which is comparatively more realistic. Let P be any point in space with coordinate xⁱ with respect to axes OXⁱ, O as origin where i = 1, 2, 3. At time t, $v^k = v^k$ (xⁱ, t) be velocity of blood, p = p (xⁱ, t) thermodynamical pressure and $\rho = \rho$ (xⁱ, t) density. Since blood vessels are cylindrical the governing equations have to transform into cylindrical co-ordinates system

3. Formulation

According to Sherman I.W. and Sherman V.G. blood is mixed fluid. There are two phases in the blood, one is plasma and other is blood cells. The blood cells are enclosed with a semi–permeable membrane whose density is greater than that of plasma. These blood cells are uniformly distributed in plasma.

3.1 Equation of continuity for two phase blood flow

According to Upadhyay V. the flow of blood is affected by the presence of blood cells. This effect is directly proportional to the volume occupied by blood cells. Let the volume portion covered by blood cells in unit volume be X, where $X = \frac{H}{100}$ and H is hematocrit the volume percentage of blood cells. Then the volume portion covered by plasma will be 1-X. If the mass ratio of blood cells to plasma is r then

$$r = \frac{X\rho_c}{(1-X)\rho_p} \tag{3.1}$$

Where ρ_c and ρ_p are densities of blood cells and plasma respectively. The both phase the blood cells and plasma move with common velocity. Campbell and Pitcher have presented a model for this condition. Equation of continuity for two phase according to principle of conservation of mass defined by J.N. and Gupta R. C. as follows

$$\frac{\partial (X\rho_c)}{\partial t} + (X\rho_c v^i)_{,i} = 0 \tag{3.2}$$

And

$$\frac{\partial (1-X)\rho_p}{\partial t} + \left((1-X)\rho_p v^i\right)_{,i} = 0$$
(3.3)

Where v velocity of mixture of two is phase blood cells and plasma, $(X\rho_c v^i)_{,i}$ is covariant derivative of $(X\rho_c v^i)$ with respect to x^i and $((1-X)\rho_p v^i)_{,i}$ is covariant derivative of $((1-X)\rho_p v^i)$ with respect to x^i .

If ρ_m be uniform density of blood then

 $\frac{1+r}{\rho_m} = \frac{r}{\rho_c} + \frac{1}{\rho_p}$

Where

$$\rho_m = X\rho_c + (1-X)\rho_p \tag{3.4}$$

Combined equation (3.2) and (3.3) and using (3.4) we get

$$\frac{\partial \rho_m}{\partial t} + (\rho_m v^i)_{,i} = 0 \tag{3.5}$$

3.2 Equation of motion for two phase blood flow-

According to Ruch T.C. and H.D. the hydro dynamical pressure p between two phases of can be supposed to be uniform because the both phases are always in equilibrium state in blood (1973). According to principle of conservation of momentum the equation of motion of two phase blood cells and plasma

$$X\rho_c \frac{\partial v^i}{\partial t} + (X\rho_c v^j)v^i_{,j} = -Xp_{,j}g^{ij} + X\eta_c \left(g^{jk}v^i_{,k}\right)_{,j}$$
(3.6)

And
$$(1-X)\rho_p \frac{\partial v^i}{\partial t} + \{(1-X)\rho_p v^j\} v^i_{,j} = -(1-X)p_{,j}g^{ij} + (1-X)\eta_p (g^{jk}v^i_{,k})_{,j}$$
(3.7)

Now adding (3.6) and (3.7) and using (3.4) then equation of motion for blood flow will be

$$\rho_m \frac{\partial v^i}{\partial t} + (\rho_m v^j) v^i_{,j} = -p_{,j} g^{ij} + \eta_m \left(g^{jk} v^i_{,k} \right)_{,j}$$
(3.8)

Where $\eta_m = X\eta_c + (1 - X)\eta_p$ is the viscosity coefficient of blood as a mixture of two phases. As velocity of blood flow decreases, the viscosity of blood increases. Since the venules are remote from heart therefore velocity of blood decreases. The Herschel Bulkley law hold good on two phase blood flow through the venules and whose constitutive equation as follows.

$$T' = \eta_m e^n + T_p (T' > T_p)$$
 and $e = 0 (T' < T_p)$ where T_p is yield stress.

When strain rate e=0 ($T' < T_p$) a core region is formed which flow just like a plug. Let radius of plug is r_p and the stress acting on the surface of plug will be T_{p} .

Equation of force acting on the plug,

$$P \pi r_p^2 = T_p 2 \pi r_p \text{ or, } r_p = 2 \frac{r_p}{P}$$

The constitutive equation for rest part of blood vessel is

 $T'_{e} = \eta_{m}e^{n} + T_{p}$ or, $T' - T_{p} = \eta_{m}e^{n} = T_{e}$ where T_{e} is effective stress whose generalized form will be $T^{ij} = -pg^{ij} + T_{e}^{ij}$ where $T_{e}^{ij} = \eta_m \ (e^{ij})^n, \ e^{ij} = g^{jk} v_k^i$

Equation of continuity $-\frac{1}{\sqrt{g}}(\sqrt{g}v^i)_{,i}=0$ (3.9)

Equation of motion - $\rho_m \frac{\partial v^i}{\partial t} + \rho_m v^j v^i_{,j} = -T^{ij}_{e,j}$ (3.10)

Where all the symbols have their usual meaning

Newton Raphson Method: The General Newton Raphson Method Formula is

$$x_{n+1} = x_n - \frac{f(x_n)}{f'(x_n)}$$
(3.11)

The above formula is repeated until a sufficiently precise value is obtained.

4. Solution

Let $x^1 = r$, $x^2 = \Theta$ and $x^3 = z$ be cylindrical co-ordinates and square length of small element $ds^2 = dr^2 + r^2 d\Theta^2 + dz^2$ Christoffel's symbols of first and second kind are given below.

$$[i j, k] = \frac{1}{2} \left[\frac{\partial g_{jk}}{\partial x^i} + \frac{\partial g_{ik}}{\partial x^j} - \frac{\partial g_{ij}}{\partial x^k} \right] \text{ and } \begin{cases} k \\ ij \end{cases} = g^{k\alpha} [ij, \alpha]$$

 $[g_{ii}]$ be matrix of metric tensor and $[g^{ij}]$ be matrix of conjugate matrix tensor where

$$\begin{bmatrix} g_{ij} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & r^2 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} g^{ij} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \frac{1}{r^2} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Metric elements $g_{rr} = 1$, $g_{\theta\theta} = r^2$, $g_{zz} = 1$

Or
$$g_{11} = 1$$
, $g_{22} = r^2$, $g_{33} = 1$

Christoffel's symbols of second kind for cylindrical co-ordinates Except of these all are zero.

Physical components

Since $\sqrt{g_{11}} v^1 = v_r$ or, $v_r = v^1$

 $\sqrt{g_{22}} v^2 = v_\theta$ or, $v_\theta = rv^2$

and
$$\sqrt{g_{33}}v^3 = v_z \, or, v_z = v^3$$

Matrix of physical components of shearing stress tensor

$$T^{\prime i j} = \eta_m \left(e^{i j} \right)^n = \eta_m \left(g^{j k} v^i_{,k} + g^{i k} v^j_{,k} \right)^n$$

$$T^{\prime i j} = \begin{bmatrix} 0 & 0 & \eta_m \left(\frac{d v}{d r} \right)^n \\ 0 & 0 & 0 \\ \eta_m \left(\frac{d v}{d r} \right)^n & 0 & 0 \end{bmatrix}$$
(4.1)

The covariant derivative of T'^{ij}

$$T'_{j}^{ij} = \frac{1}{\sqrt{g}} \frac{\partial}{\partial} \left(\sqrt{g} T^{ij} + \begin{cases} i \\ j \\ k \end{cases} \right)$$
(4.2)

According the above facts, the governing tensorial equation can be transformed into cylindrical form which is as follow

The equation of continuity $\frac{\partial v}{\partial z} = 0$ (4.3)

The equation of motion

$$\mathbf{r} - \text{Component} - \frac{\partial p}{\partial r} = 0 \tag{4.4}$$

 Θ - Component 0 = 0 (4.5)

z- Component

$$-\frac{\partial p}{\partial z} + \frac{\eta_m}{r} \frac{\partial}{\partial r} \left[r \left(\frac{\partial v_z}{\partial z} \right)^n \right]$$
(4.6)

Here this fact has been taken in view that the blood flow is axially symmetric in venules concerned i.e. $v_{\theta} = 0$ and v_r , v_z and p do not depend upon Θ and also blood flow radialy

$$\frac{\partial p}{\partial t} = \frac{\partial v_r}{\partial t} = \frac{\partial v_{\theta}}{\partial t} = \frac{\partial v_z}{\partial t} = 0$$

From (4.3)
$$v_z = v(r)$$
 (4.7)

since v_z does not depend upon θ

From equation (4.4)
$$p = p(z)$$

$$(4.8)$$

Because p does not depend upon, Θ using equation (4.7) & (4.8) in (4.6) then

$$-\frac{dp}{dz} + \frac{\eta_m}{r}\frac{d}{dr}\left[r\left(\frac{dv}{dr}\right)^n\right] = 0$$

The pressure gradient $-\frac{dp}{dz} = P$ of blood flow in venules remote from heart can be supposed to be constant, therefore equation (4.9) takes the following form

$$\frac{d}{dr}\left[\mathbf{r}\left(\frac{dv}{dr}\right)^{n} = -\frac{Pr}{\eta_{m}}\right]$$
(4.10)

On integrating (4.10), we get

$$r\left(\frac{dv}{dr}\right)^n = -\frac{pr^2}{2\eta_m} + C (4.11)$$

Since the velocity of blood flow on the axis of the cylindrical venules is maximum and constant so we apply the boundary condition at r=0, $v = v_0$ (constant) on equation (4.11) we get C=0 then equation (4.11) takes the following form

$$r\left(\frac{dv}{dr}\right)^n = -\frac{p r^2}{2\eta_m} \text{ or, } -\frac{dv}{dr} = \left(\frac{pr}{2\eta_m}\right)^{1/n}$$
(4.12)

Replace r from r - r_p for non plug region

$$\frac{dv}{dr} = -\left(\frac{P}{2\eta_m}\right)^{\frac{1}{n}} \left(r - r_p\right)^{\frac{1}{n}}$$
(4.13)

Integrating equation (4.13), we get

$$v = -\left(\frac{p}{2\eta_m}\right)^{1/n} \frac{(r-r_p)^{\frac{1}{n}+1}}{\frac{1}{n}+1} + \mathbf{B}$$
(4.14)

Apply no slip boundary condition v = 0 at r = R in equation (4.14), we get

$$B = \left(\frac{p}{2\eta_m}\right)^{1/n} \frac{(R - r_p)^{\frac{1}{n} + 1}}{\frac{1}{n} + 1}$$
(4.15)

Using (4.15) in equation 4.14) then

$$v = \left(\frac{p}{2\eta_m}\right)^{\frac{1}{n}} \frac{n}{n+1} \left[(R - r_p)^{\frac{1}{n+1}} - (r - r_p)^{\frac{1}{n+1}} \right]$$
(4.16)

This is velocity of blood flow in venules

Putting $r = r_p$ we get the velocity of plug flow as follows

$$\nu_p = \frac{n}{n+1} \left(\frac{P}{2\eta_m}\right)^{1/n} \left(R - r_p\right)^{\frac{1}{n+1}}$$
(4.17)

Where value of r_p taken from equation of motion

5. Result

The flow flux of two phased blood flow in venule

$$Q = \int_{0}^{r_{p}} 2\pi r v_{p} dr + \int_{r_{p}}^{R} 2\pi r v dr$$

= $\int_{0}^{r_{p}} 2\pi r \frac{n}{n+1} \left(\frac{P}{3\eta_{m}}\right)^{\frac{1}{n}} (R - r_{p})^{\frac{1}{n+1}} dr + \int_{r_{p}}^{R} 2\pi r \frac{n}{n+1} \left(\frac{P}{2\eta_{m}}\right)^{\frac{1}{n}} [(R - r_{p})^{\frac{1}{n+1}} - (r - r_{p})^{\frac{1}{n+1}}] dr$
(Using (4.16) & (4.17)

$$Q = \frac{2\pi n}{n+1} \left(\frac{p}{3\eta_m}\right)^{\frac{1}{n}} \left(R - r_p\right)^{\frac{1}{n}+1} \left[\frac{r^2}{2}\right]_0^{r_p} + \frac{2\pi n}{n+1} \left(\frac{p}{2\eta_m}\right)^{\frac{1}{n}} \left[\frac{r^2}{2} \left(R - r_p\right)^{\frac{1}{n}+1} - \frac{r(r - r_p)^{\frac{1}{n}}}{\frac{1}{n+2}} + \frac{(r - r_p)^{\frac{1}{n}+1}}{\left(\frac{1}{n}+2\right)\left(\frac{1}{n}+3\right)}\right]_{r_p}^{R}$$

(5.2)

$$\mathbf{Q} = \frac{\pi \mathbf{n}}{\mathbf{n}+1} \left(\frac{P}{3\eta_m}\right)^{\frac{1}{n}} \left(R\right)^{\frac{1}{n}+3} \left[\frac{r_p^2}{R^2} \left(1-\frac{r_p}{R}\right)^{\frac{1}{n}+1} + \left(1+\frac{r_p}{R}\right) \left(1-\frac{r_p}{R}\right)^{\frac{1}{n}+2} + \frac{2\left(1-\frac{r_p}{R}\right)^{\frac{1}{n}+3}}{\left(\frac{1}{n}+2\right)\left(\frac{1}{n}+3\right)} - \frac{2\left(1-\frac{r_p}{R}\right)^{\frac{1}{n}+2}}{\left(\frac{1}{n}+2\right)}\right]$$

Observations: Hemoglobin and blood pressure is taken from Anurag Nursing Home Banda by Dr. Anurag Srivstava Patient Name – Gyanchandra, Age – 43 years / male Annual No.- 0645/2022 Diagnosis – Falsiparum Malaria

Table 1: 1	Hematocrit and	Cli.	Blood	Pressure	Drop
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Sl. No.	Date	HB (Hemoglobin) gm/dl	BP (Blood Pressure) mmhg	Hematocrit (3×HB)	BP (In Pascal)	Blood Pressure Drop
1	22.10.2022	11.7	135/92	35.1	17998.2/12265.44	-3044.14
2	24.10.2022	11.4	142/94	34.2	18931.44/12532.08	-3140.43
3	26.10.2022	10.7	140/90	32.1	18664.8/11998.8	-3036.73
4	28.10.2022	10.1	148/96	30.3	19731.36/12798.72	-3229.31
5	30.10.2022	10.4	138/90	31.2	18398.16/11998.8	-3021.92

Blood Pressure Drop in $=\frac{2}{3} \left[\frac{\frac{S+D}{2}+D}{3}\right] - \left[\frac{\frac{S+D}{2}+D}{3}\right]$

 $Q = 1000 \text{ ml/min} = 0.01666 \text{ litre/sec} = 0.0000167 \text{ m}^3\text{/sec}$

$$R = 1, r_p = 1/3$$

According to Gustafason Daniel R, (1980) [11]

 $\eta_p = 0.0015$ (Pascal-Sec)

According to Glenn Elert (2010)

 $\eta_m = 0.035$ (Pascal –Sec)

Average length of terminal hepatic venule = 0.15cm = 15×10^{-4} meter

Average systolic pressure S= 18744.792 and average diastolic pressure D= 12318.768

H= 32.58, Blood pressure drop = 3094.50533 (Pascal Sec.)

Since $\eta_m = \eta_c X + \eta_p$ (1-X) where $X = \frac{H}{100}$

Substituting the values of η_m , η_p , and H in above relation we get $\eta_c = 0.104323818$ again from above relation

 $\eta_m = 0.001028238187 \mathrm{H} + 0.0015$

Substituting the value of r_p and R in equation (28) we get

$$Q = \frac{2\pi}{27} \left(\frac{P}{3\eta_m}\right)^{\frac{1}{n}} \left[\frac{26n^3 + 33 n^2 + 9n}{6n^3 + 11 n^2 + 6n + 1}\right]$$

Or, $\frac{27Q}{2\pi} = \left(\frac{P}{3\eta_m}\right)^{\frac{1}{n}} \left[\frac{26n^3 + 33 n^2 + 9n}{6n^3 + 11 n^2 + 6n + 1}\right]$
Or, $\left(\frac{P}{3\eta_m}\right) = \left(\frac{27Q}{2\pi A}\right)^n$ where $A = \frac{26n^3 + 33 n^2 + 9n}{6n^3 + 11 n^2 + 6n + 1}$
Or, $P = \left(\frac{27Q}{2\pi A}\right)^n 3\eta_m$
Since $P = -\frac{dp}{dz}$ or, $dp = -Pdz$
 $p_f - p_i = \left(\frac{27Q}{2\pi A}\right)^n 3\eta_m (z_f - z_i)$

Where $p_f - p_i$ pressure drop and $z_f - z_i$ = length of hepatic venule Substituting the values of Q, $p_f - p_i$, and η_m in above equation and solve by numerical method we get n = -1.4445, and again

$$p_f - p_i = 3\eta_m (z - z_i) \left(\frac{27Q}{2\pi A}\right)^n$$

Substituting the value of $3\eta_m$, Q and n we get

$p_f - p_i = 90.975056 \text{ H} + 132.714954$

This is relation between hematocrit and blood pressure drop

Table 2: Hematocrit and Modulated Blood Pressure Drop

Date	22.10.2022	24.10.2022	26.10.2022	28.10.2022	30.10.2022
Hematocrit	35.1	34.2	32.1	30.3	31.2
Blood pressure drop	3325.94	3244.06	3053.01	2889.26	2971.17



Graph 1: Hematocrit Vs Modulated Blood Pressure Drop



Graph 2: Hematocrit Vs (Modulated and Cli.) Blood Pressure drop

6. Observation of graph

The graph-1 shows that fluctuations of blood pressure drop with respect to hematocrit of five different dates. We observed minimum blood pressure drop 2889.26 on dated 28/10/2022 and maximum value obtains 3325.94 on dated 22/10/2022. At the hematocrit value from 35.1 to 30.3 via 34.2 and 32.1 the blood pressure drop straightly decreases on dated from 22/10/2022 to 28/10/2022 via 24/10/2022 and

26/10/2022 and hematocrit value from 30.3 to 31.2 the blood pressure drop straightly increases on dated from 28/10/2022 to 30/10/2022. The graph-2 shows that the comparative study of two graph (i) Graph between hematocrit Vs clinical blood pressure drop (ii) Graph between hematocrit Vs mathematically modulated blood pressure drop.

7. Conclusion

(5.3)

In graph-1 the slope of straight line is the absolute value and when hematocrit increases the blood pressure drop also increases and when hematocrit decreases the blood pressure drop also decreases. When the trend of straight line decreasing sense then medicine dose slowly increases, when steepness of curve low then we can give high dose of medicine and when trend of straight line increases sense then we suggest normal dose of medicine. A comparative study of both graphs shows that nearly have the same character.

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